



GOALS-JWST: Pulling Back the Curtain on the AGN and Star Formation in VV 114

J. Rich¹, S. Aalto², A. S. Evans^{3,4}, V. Charmandaris^{5,6,7}, G. C. Privon^{3,8}, T. Lai⁹, H. Inami¹⁰, S. Linden¹¹, L. Armus⁹, T. Diaz-Santos^{6,7}, P. Appleton⁹, L. Barcos-Muñoz⁴, T. Böker¹², K. L. Larson¹³, D. R. Law¹⁴, M. A. Malkan¹⁵, A. M. Medling^{16,17}, Y. Song^{3,4}, V. U¹⁸, P. van der Werf¹⁹, T. Bohn¹⁰, M. J. I. Brown²⁰, L. Finnerty¹⁵, C. Hayward²¹, J. Howell⁹, K. Iwasawa^{22,23}, F. Kemper^{23,24,25}, J. Marshall²⁶, J. M. Mazzarella⁹, J. McKinney²⁷, F. Muller-Sanchez²⁸, E. J. Murphy²⁹, D. Sanders³⁰, B. T. Soifer³¹, S. Stierwalt³², and J. Surace⁹

¹The Observatories of the Carnegie Institution for Science, 813 Santa Barbara Street, Pasadena, CA, 91101, USA

²Department of Space, Earth and Environment, Chalmers University of Technology, SE-412 96 Gothenburg, Sweden

³National Radio Astronomy Observatory, 520 Edgemont Rd, Charlottesville, VA, 22903, USA

⁴Department of Astronomy, University of Virginia, 530 McCormick Road, Charlottesville, VA, 22903, USA

⁵Department of Physics, University of Crete, Heraklion, 71003, Greece

⁶Institute of Astrophysics, Foundation for Research and Technology-Hellas (FORTH), Heraklion, 70013, Greece

⁷School of Sciences, European University Cyprus, Diogenes street, Engomi, 1516 Nicosia, Cyprus

⁸Department of Astronomy, University of Florida, P.O. Box 112055, Gainesville, FL, 32611, USA

⁹IPAC, California Institute of Technology, 1200 E. California Blvd., Pasadena, CA, 91125, USA

¹⁰Hiroshima Astrophysical Science Center, Hiroshima University, 1-3-1 Kagamiyama, Higashi-Hiroshima, Hiroshima 739-8526, Japan

¹¹Department of Astronomy, University of Massachusetts at Amherst, Amherst, MA, 01003, USA

¹²European Space Agency, Space Telescope Science Institute, Baltimore, MD, 21218, USA

¹³AURA for the European Space Agency (ESA), Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD, 21218, USA

¹⁴Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD, 21218, USA

¹⁵Department of Physics & Astronomy, 430 Portola Plaza, University of California, Los Angeles, CA, 90095, USA

¹⁶Department of Physics & Astronomy and Ritter Astrophysical Research Center, University of Toledo, Toledo, OH, 43606, USA

¹⁷ARC Centre of Excellence for All Sky Astrophysics in 3 Dimensions (ASTRO 3D); Australia

¹⁸Department of Physics and Astronomy, 4129 Frederick Reines Hall, University of California, Irvine, CA, 92697, USA

¹⁹Leiden Observatory, Leiden University, PO Box 9513, 2300 RA Leiden, The Netherlands

²⁰School of Physics and Astronomy, Monash University, Clayton, VIC 3800, Australia

²¹Center for Computational Astrophysics, Flatiron Institute, 162 Fifth Avenue, New York, NY, 10010, USA

²²Institut de Ciències del Cosmos (ICCUB), Universitat de Barcelona (IEEC-UB), Martí i Franquès, 1, E-08028 Barcelona, Spain

²³ICREA, Pg. Lluís Companys 23, E-08010 Barcelona, Spain

²⁴Institut de Ciències de l'Espai (ICE, CSIC), Can Magrans, s/n, E-08193 Bellaterra, Barcelona, Spain

²⁵Institut d'Estudis Espacials de Catalunya (IEEC), E-08034 Barcelona, Spain

²⁶Glendale Community College, 1500 N. Verdugo Rd., Glendale, CA, 91208, USA

²⁷Department of Astronomy, University of Massachusetts, Amherst, MA, 01003, USA

²⁸Department of Physics and Materials Science, The University of Memphis, 3720 Alumni Avenue, Memphis, TN, 38152, USA

²⁹National Radio Astronomy Observatory, 520 Edgemont Road, Charlottesville, VA, 22903, USA

³⁰Institute for Astronomy, University of Hawaii, 2680 Woodlawn Drive, Honolulu, HI, 96822, USA

³¹Division of Physics, Mathematics and Astronomy, California Institute of Technology, 1200 E. California Blvd., Pasadena, CA, 91125, USA

³²Physics Department, 1600 Campus Road, Occidental College, Los Angeles, CA, 90041, USA

Received 2022 November 22; revised 2022 December 20; accepted 2023 January 4; published 2023 February 22

Abstract

We present results from the James Webb Space Telescope Director's Discretionary Time Early Release Science program 1328 targeting the nearby, luminous infrared galaxy, VV 114. We use the MIRI and NIRSpec instruments to obtain integral-field spectroscopy of the heavily obscured eastern nucleus (V114E) and surrounding regions. The spatially resolved, high-resolution spectra reveal the physical conditions in the gas and dust over a projected area of 2–3 kpc that includes the two brightest IR sources, the NE and SW cores. Our observations show for the first time spectroscopic evidence that the SW core hosts an active galactic nucleus as evidenced by its very low 6.2 μm and 3.3 μm polycyclic aromatic hydrocarbon equivalent widths (0.12 and 0.017 μm , respectively) and mid- and near-IR colors. Our observations of the NE core show signs of deeply embedded star formation including absorption features due to aliphatic hydrocarbons, large quantities of amorphous silicates, as well as HCN due to cool gas along the line of sight. We detect elevated [Fe II]/P α consistent with extended shocks coincident with enhanced emission from warm H₂, far from the IR-bright cores and clumps. We also identify broadening and multiple kinematic components in both H₂ and fine structure lines caused by outflows and previously identified tidal features.

Unified Astronomy Thesaurus concepts: [Interacting galaxies \(802\)](#); [Galaxy nuclei \(609\)](#); [Luminous IR galaxies \(946\)](#); [Active galactic nuclei \(16\)](#)

Supporting material: machine-readable table

1. Introduction

VV 114 (Arp236, IC1623) is an interacting system undergoing vigorous starburst activity. With an IR luminosity of $L_{\text{IR}} \sim 4.5 \times 10^{11} L_{\odot}$, and a distance of 80 Mpc, it is one of the

brightest objects in the IRAS Bright Galaxy Sample (Soifer et al. 1987). It appears to be an early-stage merger of two galaxies that are aligned east–west with a projected nuclear separation of ~ 8 kpc, designated in the literature as VV 114E and VV 114W, respectively. At optical wavelengths, VV 114 shows a highly disturbed morphology with very faint tidal tails extending over 25 kpc from the center (Arp 1966). The western component, VV 114W, is more extended than the eastern one, and dominates the emission at short wavelengths. Much of the mid-IR emission is diffuse and extended over several kpc with some indication of an active galactic nucleus (AGN) based on the mid-IR color of the more compact nuclear region in VV 114E (Le Floch et al. 2002). Atacama Large Millimeter/submillimeter Array (ALMA) observations show abundant cold, dense gas (traced by e.g., CO, HCO, HCN), evidence for shocked gas in the overlap region between the two galaxies (traced by methanol), a molecular outflow, and a possible buried AGN in VV 114E (Iono et al. 2013; Saito et al. 2015, 2017, 2018). The majority of the IR emission and by extension total energy output of the system is dominated by VV 114E and the extreme UV/optical-to-IR ratios of VV 114 make it a more plausible analog to high- z IR luminous mergers (Charmandaris et al. 2004; Howell et al. 2010).

Mid-IR spectra centered on VV 114E taken with Spitzer IRS ($\sim 10'' \times 36''$) show a moderately strong $9.7 \mu\text{m}$ silicate absorption ($S_{9.7\mu\text{m}} = -0.98$), a $6.2 \mu\text{m}$ polycyclic aromatic hydrocarbon (PAH) equivalent width (EW) intermediate to luminous infrared galaxies (LIRGs) dominated by star formation or AGN ($\text{EW}_{6.2\mu\text{m}} = 0.30 \mu\text{m}$), and a ratio of H_2 to PAH luminosity slightly above values associated with photodissociation region emission (Guillard et al. 2012a; Stierwalt et al. 2013, 2014). No coronal lines (e.g., [Ne V]) were detected in the Spitzer spectra, and the fine structure line flux ratios were consistent with LIRGs primarily dominated by star formation with some composite AGN/starburst activity (Inami et al. 2013). Observations with Chandra, XMM, and NuSTAR of VV 114E indicate that the X-ray emission appears to be generated primarily through star formation, with little to no X-ray emission coming from an AGN (Grimes et al. 2006; Garofali et al. 2020; Ricci et al. 2021). Visible wavelength integral-field spectroscopy indicates a mix of star formation and shock emission, the latter indicated by elevated emission line ratios and line widths across both galaxies (Rich et al. 2011, 2015). Finally, in a companion paper to the present study, Evans et al. (2022) propose the presence of a reddened starburst and an AGN in the bright NE and SW cores, respectively, of VV 114E based on broadband JWST mid-IR colors.

In this paper we present combined JWST MIRI/NIRSpec integral-field spectroscopic observations of the nucleus of VV 114E and surrounding regions taken as part of the Early Release Science (ERS) program 1328 (Co-PIs: L. Armus and A. Evans). The data allow us to resolve the properties of the two brightest sources, the NE and SW cores, as well as star clusters and diffuse emission surrounding the eastern nucleus.

Throughout the paper we adopt a cosmology of $H_o = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.28$, and $\Omega_\Lambda = 0.72$. The redshift of VV 114 ($z = 0.0202$) corresponds to an angular scale of $400 \text{ pc}/1''$ (Wright 2006).

2. Observations and Data Reduction

VV 114 was observed with MIRI (Rieke et al. 2015; Labiano et al. 2021) on 2022 July 2 (MIRI imaging; Bouchet et al. 2015), on 2022 July 5 (MRS spectroscopy; Wells et al. 2015), and by NIRSpec (Jakobsen et al. 2022) on 2022 July 19 (IFU spectroscopy; Böker et al. 2022). We include MIRI imaging in this paper to indicate the locations of our spectral extraction apertures. The full analysis of these imaging data products are described by Evans et al. (2022).

2.1. MIRI MRS Data

The MIRI MRS observations include three grating settings (SHORT, MEDIUM, and LONG) in order to cover the entire wavelength range accessible to the four IFU channels ($4.9\text{--}27.9 \mu\text{m}$). A four-point dither pattern was employed to recover the extended emission and avoid saturation, with separate off-source observations for background subtraction. The field of view (FOV) and position angle (PA) vary by channel, but the FOV with full wavelength coverage is defined by channel 1, $\sim 5'' \times 4''$ at $\text{PA} = 255^\circ$ as shown in Figure 1.

Our procedure for reducing the MRS data of the ERS 1328 targets is described in detail by U et al. (2022), but a brief summary is given here: uncalibrated data are processed using the most recently available developmental release of the JWST Science Calibration Pipeline (Bushouse et al. 2022), version 1.8.3, using calibration reference data system (CRDS) context file `jwst_0963.pmap`. The resulting data products generated by the pipeline are 12 background-subtracted, fringe-corrected, wavelength- and flux-calibrated data cubes—one for every combination of channel and grating setting combining the four dither pointings. These cubes are used to generate 1D spectra of regions of interest and 2D maps of the properties of particular emission line features.

We extract spectra from 10 regions of interest within the channel 1 FOV to investigate IR-bright sources resolved in the MIRI imaging as well as diffuse emission surrounding those sources and in the eastern nucleus (Figure 1). We chose five locations coincident with the two brightest IR sources, the NE and SW cores ((a), (c)), and several bright clumps identified in the MIRI imaging data and in both the submillimeter and radio ((d), (e), and (f)) (Saito et al. 2018; Evans et al. 2022; Song et al. 2022). We also chose five regions intended to capture diffuse emission between (b) and around ((g), (h), (i), and (j)) the bright clumps coincident with some tidal and shock features previously identified using ALMA (Saito et al. 2017, 2018).

Spectra are extracted from each of the 12 data cubes using apertures with a radius of $0''.4$ (160 pc). This results in 12 1D spectra for each extracted region, with some overlap in adjacent wavelength regions. The 12 individual bands for each spectra have slight offsets in flux (a few percent) which are multiplicatively scaled, trimmed, and stitched in order to create continuously smooth 1D spectra over the full MIRI wavelength range. This process begins by using the overlapping wavelength region to scale the Channel 4 MEDIUM spectrum to the longest wavelength Channel 4 LONG spectrum ($\sim 23\text{--}25 \mu\text{m}$), trimming the overlapping values from the noisier spectrum of the two, and continuing the process to channels at shorter and shorter wavelengths. Finally, a wavelength-dependent aperture correction is applied to each spectrum (U et al. 2022).

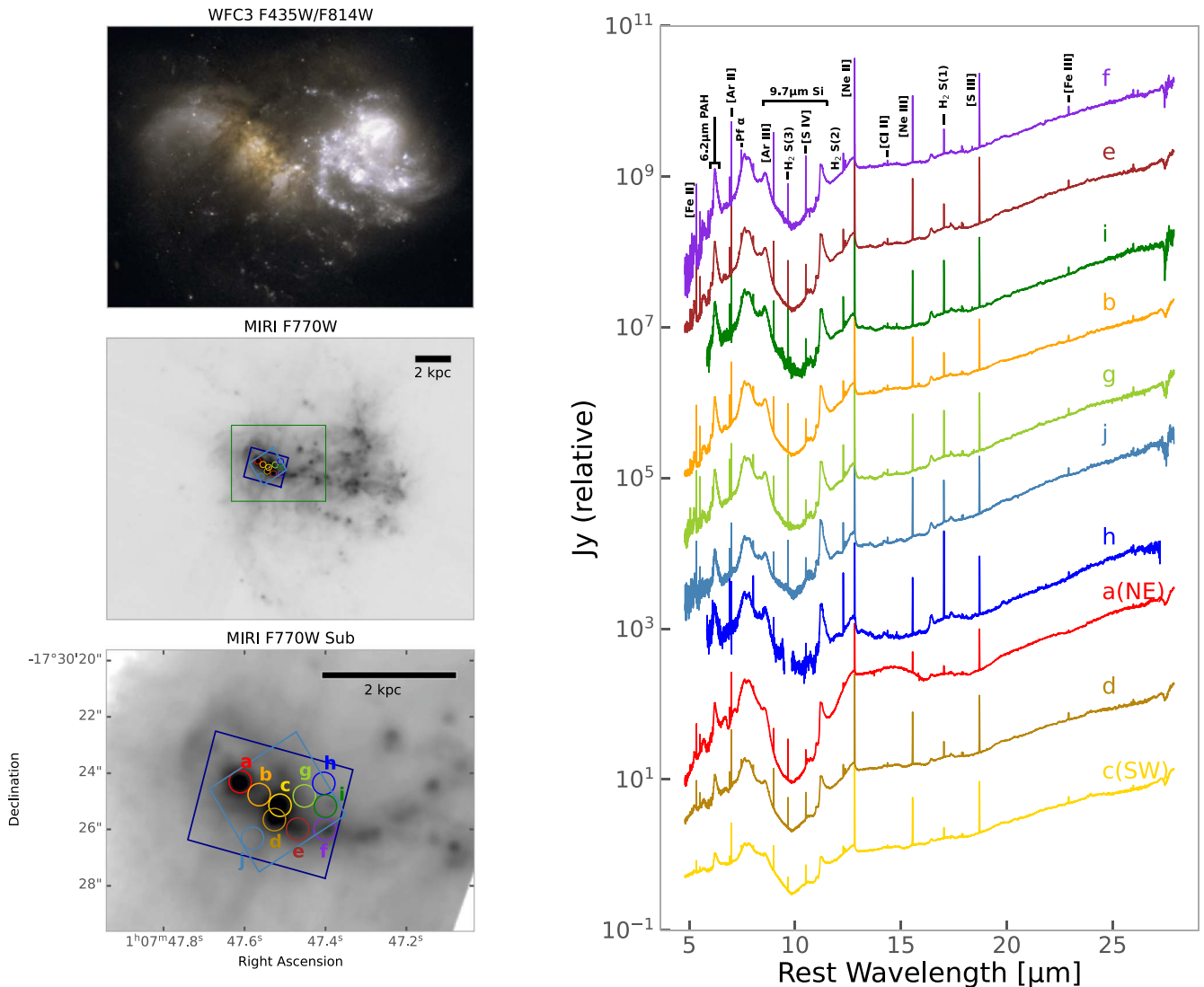


Figure 1. Images and spectra of VV 114. Top left: Hubble Space Telescope (HST) 435W/814W color image. Middle left: JWST F770W image with the same FOV, where the green box corresponds to the MIRI SUB128 subarray FOV, as shown in the bottom left panel. Bottom left: F770W SUB128 image with MIRI MRS channel 1 FOV (dark blue box), NIRSpect FOV (light blue box), and our ten $0''.4$ radius extraction apertures. Right: full MIRI spectra of the 10 regions marked in the F770W subarray image. Spectra are shown at rest-frame wavelengths assuming a systemic velocity of 6056 km s^{-1} and sorted from top to bottom in order of decreasing $6.2 \mu\text{m}$ PAH EW.

2.2. NIRSpect IFU Data

NIRSpect IFU observations were taken with three grating and filter combinations: G140H/F100LP, G235H/F170LP, and G395H/F290LP to cover the spectral range from 0.97 to $5.3 \mu\text{m}$. Calculations in this paper were made using the wavelength region covered by the G235H/F170LP and G395H/F290LP settings, $\sim 1.7\text{--}5.3 \mu\text{m}$. This wavelength range allows us to measure the AGN-sensitive $3.3 \mu\text{m}$ PAH feature. Again a four-point dither pattern was employed to sample the PSF completely and to avoid saturation, with an additional “Leakcal” image taken for each grating setting. The FOV of the combined dither pattern is $\sim 3''.6 \times 3''.8$ centered on the SW Core at a PA of $\sim 32^\circ$ (Figure 1).

We reduce the NIRSpect data in a similar fashion to the MRS data, using CRDS context file `jwst_1009.pmap`. Uncalibrated four-point dither pattern science images and Leakcal images, one for each of the two NIRSpect chips, are downloaded using MAST resulting in 10 image files. These are first put through Stage 1 processing with the `Detector1` pipeline, which

applies detector-level calibrations and produces count rate files calculated from the nondestructive “ramp” readouts. These rate files are then processed using the `Spec2` pipeline, which applies physical corrections and flux and wavelength calibrations. At this step in the overall pipeline, the Leakcal files are also used to correct for any stray light that may fall on the detector due to failed open MSA shutters. Finally we run the `Spec3` pipeline step, which produces a final combined data cube sampled with $0''.1$ spaxels.

For our analysis in this paper, we extract from the final data cube spectra from the two brightest IR sources, the NE and SW cores ((a) and (c), respectively). We matched our apertures to the MIRI MRS extraction radius of $0''.4$ centered at the same two locations. For the G395H/F290LP setting this produces flux- and wavelength-calibrated 1D spectra covering $2.87\text{--}5.27 \mu\text{m}$, with a gap in coverage from 4.06 to $4.18 \mu\text{m}$ in the middle of the spectrum due to the gap between the two NIRSpect chips. For G235H/F170LP the range is $1.66\text{--}3.17 \mu\text{m}$ with a gap from 2.40 to $2.45 \mu\text{m}$. We use the

Table 1
Spectral Feature Strengths, Fluxes, and FWHM

Region ID	EW _{3.3μm}	EW _{6.2μm}	S _{9.7}	[Fe II] 5.34	FWHM	Pfund α	FWHM
a (NE Core)	0.121 \pm 0.001	0.264 \pm 0.018	-2.45 \pm 0.03	7.41 \pm 0.22	194 \pm 13	1.56 \pm 0.62	142 \pm 37
b		0.514 \pm 0.017	-1.13 \pm 0.03	11.6 \pm 0.29	186 \pm 12	2.83 \pm 0.58	148 \pm 21
c (SW Core)	0.017 \pm 0.001	0.106 \pm 0.002	-1.06 \pm 0.01	7.49 \pm 0.40	178 \pm 14	4.76 \pm 0.92	154 \pm 20
d		0.199 \pm 0.015	-1.44 \pm 0.02	9.87 \pm 0.38	166 \pm 13	4.92 \pm 0.88	135 \pm 18
e		0.652 \pm 0.015	-1.17 \pm 0.04	7.64 \pm 0.19	185 \pm 12	2.23 \pm 0.37	152 \pm 18
f		0.720 \pm 0.062	-0.90 \pm 0.04	3.09 \pm 0.13	153 \pm 13	2.55 \pm 0.37	120 \pm 13
g		0.491 \pm 0.036	-1.05 \pm 0.02	5.03 \pm 0.20	221 \pm 14	0.79 \pm 0.16	192 \pm 25
h		0.34 \pm 0.17	-0.75 \pm 0.04	2.32 \pm 0.57	147 \pm 29	0.24 \pm 0.17	160 \pm 73
i		0.56 \pm 0.14	-0.73 \pm 0.05	2.78 \pm 0.39	231 \pm 26	1.08 \pm 0.36	224 \pm 44
j		0.359 \pm 0.045	-0.82 \pm 0.01	2.30 \pm 0.14	270 \pm 17	0.24 \pm 0.07	243 \pm 43

Note. Values of the 3.3 and 6.2 μ m EWs, 9.7 μ m silicate strength, as well as the line fluxes (10^{-18} W/m²) and FWHM (km s⁻¹, corrected for instrumental broadening) of the emission line features measured in our extracted regions that were used in our analysis and discussion. This table is a subset of the total measurements, a machine-readable version of the full table is available online.

(This table is available in its entirety in machine-readable form.)

overlapping wavelength range between the NIRSpect and MIRI data to scale the G395H/F290LP spectrum to match the MIRI spectrum, and the overlapping region between the two NIRSpect settings to scale and stitch the shorter wavelength spectra (G235H/F170LP) to the longer wavelength NIRSpect spectra.

3. Results

We use the 1D aperture-extracted spectra to measure emission and absorption features that trace the physical properties of the gas and dust. Several of our apertures are centered on bright, unresolved mid-IR sources seen in the MIRI imaging observations (Evans et al. 2022), including the bright NE and SW cores ((a) and (c), respectively), a source directly SE of the SW core (d), and a deeply embedded star cluster (f) with mass $M \sim 10^6 M_{\odot}$, age $t \sim 1\text{--}2$ Myr, and extinction $A_V \sim 8$ (see Linden et al. 2022). The remaining apertures trace diffuse emission generally showing elevated EW_{6.2 μ m} and strong H₂ emission.

3.1. Emission Line Properties

We perform fits to features in each 1D MIRI spectrum using the `lmfit` package (Newville et al. 2014) with the resulting values given in Table 1. Atomic and H₂ emission lines are fit with a single Gaussian component combined with a polynomial fit to the local continuum over a range of 0.1 μ m. The resulting Gaussian parameters are used to derive the observed flux of each line. The width of the Gaussian fit is used to determine the intrinsic FWHM of each emission line by subtracting in quadrature the instrumental resolution of MRS at the observed wavelength (Labiano et al. 2021).

Several fine structure lines and H₂ lines are well detected and resolved, as well as the hydrogen recombination lines Pfund α and Humphreys α and β . The emission line ratios show variation between apertures indicative of both widespread star formation and shock excitation. We do not detect any of the high-ionization coronal lines typically found in the mid-IR spectra of AGN-dominated galaxies (e.g., [Ne V]; Genzel et al. 1998; Lutz et al. 2000; Sturm et al. 2002; Weedman et al. 2005; Armus et al. 2007) in contrast to JWST observations of NGC 7469 that show nine well-detected coronal lines excited by high-energy photons from the central AGN (Armus et al. 2023).

Emission line ratios sensitive to AGN activity (e.g., [S IV]/[Ne II] and [O IV]/[Ne II]) show values indicative of a composite of AGN and starburst activity (Inami et al. 2013). Apertures (g), (h), and (j) show elevated values of [Fe II]/Pfund α (>5) and (g) and (j) show [Fe II] FWHM values that are higher than the other apertures (200–300 km s⁻¹).

Apertures (a), (b), (c), (d), and (e) have fine structure lines ([Ar II], [Ar III], [Ne II], [Cl II], [Ne III], and [S IV]) with FWHM values ranging between 150 and 200 km s⁻¹ and show no trend with the emission line ionization potential. In the NE core (a) the H₂ FWHM is ~ 50 km s⁻¹ narrower than the fine structure lines, while in the SW core (c) the H₂ and fine structure line widths are similar. This trend is reversed in aperture (f) where the H₂ FWHM values are ~ 200 km s⁻¹ versus ~ 100 km s⁻¹ for the fine structure lines. The broadest emission line widths in our data are observed in the spectrum of aperture (i), which samples the diffuse H₂ bright gas near the western edge of our FOV close to the “overlap” region defined by Saito et al. (2017).

In order to assess the morphology and kinematics of the region observed with MIRI/MRS, we have also carried out spaxel-by-spaxel fits of the [Ne II] 12.8 μ m fine structure emission line and H₂ S(2) 12.3 μ m molecular emission line. These two emission lines are generally quite luminous and are well detected across nearly every spaxel. The lines are covered by Channel 3B (13.29–15.52 μ m) which provides a wider FOV (7" \times 6") while still maintaining relatively high spectral resolution ($R \sim 2800\text{--}3000$) with somewhat larger spaxels ($\sim 0''.2$). Fits to the two emission lines are carried out on a spaxel-by-spaxel basis using the Channel 3B sub-band data cube, with a single Gaussian component fit to each line independently, in the same manner as the aperture-extracted spectra. The resulting maps are shown in Figure 2.

The variations in the flux and FWHM of both the [Ne II] and H₂ lines agree with the values measured in our extracted apertures. The broadest FWHMs in both lines lay outside the Channel 1 FOV, and our closest aperture extractions are (i) and (j). The increase in H₂ FWHM at the NW corner of our map corresponds to a portion of the “overlap” region observed by Saito et al. (2017) with a similar increase in FWHM seen at submillimeter wavelengths.

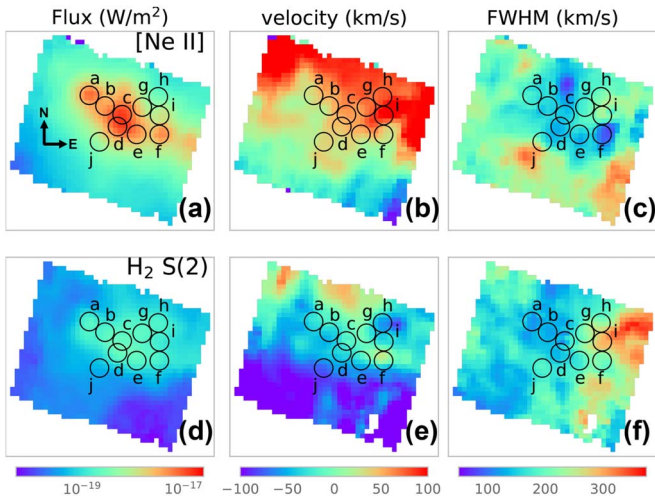


Figure 2. 2D images of the flux (a) and (d), relative velocity (b) and (e), and FWHM (c) and (f) generated from the spaxel-by-spaxel fits to the [Ne II] ($12.8 \mu\text{m}$) and $\text{H}_2 \text{S}(2)$ ($12.3 \mu\text{m}$) emission lines. The FWHMs shown are corrected for instrumental broadening as described in the text. The relative velocity map is generated by subtracting a systemic velocity of 6056 km s^{-1} from each spaxel. The broadest H_2 FWHM in the NW is just outside the Channel 1 FOV but falls within the Channel 3 FOV, and corresponds to the eastern edge of the shocked “overlap” region observed in Saito et al. (2017). The elevated FWHM in [Ne II] is near apertures that show evidence of shocks and double peaked emission line profiles ((g), (h), (i), and (j)).

3.2. PAH EW

We measure the EW of the $3.3 \mu\text{m}$ and $6.2 \mu\text{m}$ PAH emission features ($\text{EW}_{3.3 \mu\text{m}}$ and $\text{EW}_{6.2 \mu\text{m}}$, respectively) in the NIRSpec and MIRI spectra by applying the same method to both. First, portions of the spectrum adjacent to each PAH feature are used to perform a linear interpolation of the continuum. We then integrate a spline fit from 3.20 to 3.28 and 5.95 to $6.55 \mu\text{m}$ (rest frame, $\text{EW}_{3.3 \mu\text{m}}$ and $\text{EW}_{6.2 \mu\text{m}}$, respectively) to calculate the PAH feature flux. The integrated flux is divided by the continuum flux density at the wavelength of the peak of each PAH feature.

The $\text{EW}_{6.2 \mu\text{m}}$ values range from 0.11 – $0.72 \mu\text{m}$, bracketing the published Spitzer IRS value of $0.3 \mu\text{m}$ (Stierwalt et al. 2013). The NE core has an $\text{EW}_{6.2 \mu\text{m}}$ of $0.264 \mu\text{m}$ and a relatively high $\text{EW}_{3.3 \mu\text{m}}$ of $0.121 \mu\text{m}$, while the SW core has a low $\text{EW}_{6.2 \mu\text{m}}$ of $0.106 \mu\text{m}$ and a very low values of $\text{EW}_{3.3 \mu\text{m}}$ of $0.016 \mu\text{m}$, indicative of an AGN (e.g., Imanishi et al. 2010; Petric et al. 2011; see the discussion in Section 4).

The largest $\text{EW}_{6.2 \mu\text{m}}$ of $0.72 \mu\text{m}$ is measured for source (f), a very young, highly enshrouded star cluster revealed by JWST NIRCам (Linden et al. 2022). The remaining apertures have a range of values from ~ 0.20 to $0.65 \mu\text{m}$ tracing the presence of extended star formation and the influence of the AGN in the SW core.

3.3. Silicate Absorption Strength

To compare with previously published values we also calculate the apparent $9.7 \mu\text{m}$ silicate absorption feature strength in a manner consistent with Spoon et al. (2007)’s measurements of PAH-dominated sources. We assume an extrapolated power-law fit for the continuum, constrained by portions of the spectrum at $5.5 \mu\text{m}$ and $14 \mu\text{m}$ and take the natural logarithm of the ratio of the observed and interpolated continuum flux density at $9.7 \mu\text{m}$ ($s_{9.7 \mu\text{m}}$).

The majority of the $s_{9.7 \mu\text{m}}$ values measured range from -0.73 to -1.44 , bracketing the published Spitzer IRS measurement of -0.98 , with the exception of the NE core (a). Visual inspection of the NE core’s spectrum indicates a much deeper silicate absorption feature, measured to be $s_{9.7 \mu\text{m}} = -2.45$. For absorption-dominated sources Spoon et al. (2007) used a spline fit to determine a continuum. If we assume the NE core is absorption-dominated and apply the same process, the measured $s_{9.7 \mu\text{m}}$ would be stronger (-2.87). Several other features unique to the spectrum of the NE source are shown in Figure 3 including aliphatic hydrocarbon absorption and crystalline silicate absorption, both observed in ULIRGs with deep $9.7 \mu\text{m}$ silicate absorption (Spoon et al. 2007). We note that these strong absorption features are not seen in either the SW core (c) or the embedded cluster (f).

3.4. Molecular Absorption Features

The MIRI spectra allow high-resolution vibration–rotation spectroscopy of gas-phase molecules toward dust-enshrouded regions in the VV 114 system. We report the detection in absorption of the ν_2 $14.04 \mu\text{m}$ bending mode of hydrogen cyanide (HCN) and tentatively also the ν_5 $13.7 \mu\text{m}$ bending mode of acetylene (C_2H_2) toward the NE core (a). C_2H_2 is a key ingredient in the gas-phase formation of large molecules such as HC_3N , and HCN is one of the most abundant nitrogen bearing molecules in dense ($n > 1 \times 10^4 \text{ cm}^{-3}$) molecular clouds.

These lines have been previously detected by Spitzer toward dust-enshrouded young stellar objects (YSOs) (e.g., Lahuis & van Dishoeck 2000) and toward LIRGs (Lahuis et al. 2007). Lahuis et al. (2007) find HCN column densities ranging between $N(\text{HCN}) = 1$ – $12 \times 10^{16} \text{ cm}^{-2}$ and warm gas with excitation temperatures $T_{\text{ex}} = 230$ – 700 K . Owing to the higher spatial and spectral resolution, the JWST MIRI spectrum of the NE core is more complex than those found by Lahuis et al. (2007) using Spitzer.

We perform preliminary fits using the methodology of Lahuis & van Dishoeck (2000) which assumes LTE. Our preliminary results are consistent with an excitation temperature of 300 – 500 K and an $N(\text{HCN})$ of 1 – $5 \times 10^{16} \text{ cm}^{-2}$. For an HCN abundance (with respect to H_2) of 10^{-8} – 10^{-7} (e.g., Schilke et al. 1992; Lahuis & van Dishoeck 2000; Harada et al. 2013), this would be consistent with a high obscuration with $N(\text{H}_2) 10^{23}$ – 10^{24} cm^{-2} . The HCN and C_2H_2 spectra require further analysis, including multiple temperature component modeling and inclusion of non-LTE effects (V. Buiten et al. 2023, in preparation).

It is also possible that the nuclear obscuration is even higher with column densities in excess of $N(\text{H}_2) 10^{25} \text{ cm}^{-2}$ —the so called compact obscured nuclei (CONs; e.g., Aalto et al. 2015). Such objects are characterized by high millimeter and submillimeter continuum surface brightnesses and luminous emission from millimeter-wave rotational transitions of HCN within the vibrationally excited ladder (HCN-vib) (e.g., Sakamoto et al. 2010; Aalto et al. 2015, 2019; Falstad et al. 2021; Sakamoto et al. 2021). For such deeply enshrouded objects, the HCN $14 \mu\text{m}$ line, or silicate absorption, may not trace the full $N(\text{H}_2)$, but only its surface. The millimeter/submillimeter HCN-vib line (and the millimeter/submillimeter continuum) would probe deeper, revealing the full obscuring column. Falstad et al. (2021) propose that CONs have surface brightnesses of HCN-vib of $\Sigma(\text{HCN-vib}) > 1 L_{\odot} \text{ pc}^{-2}$. A

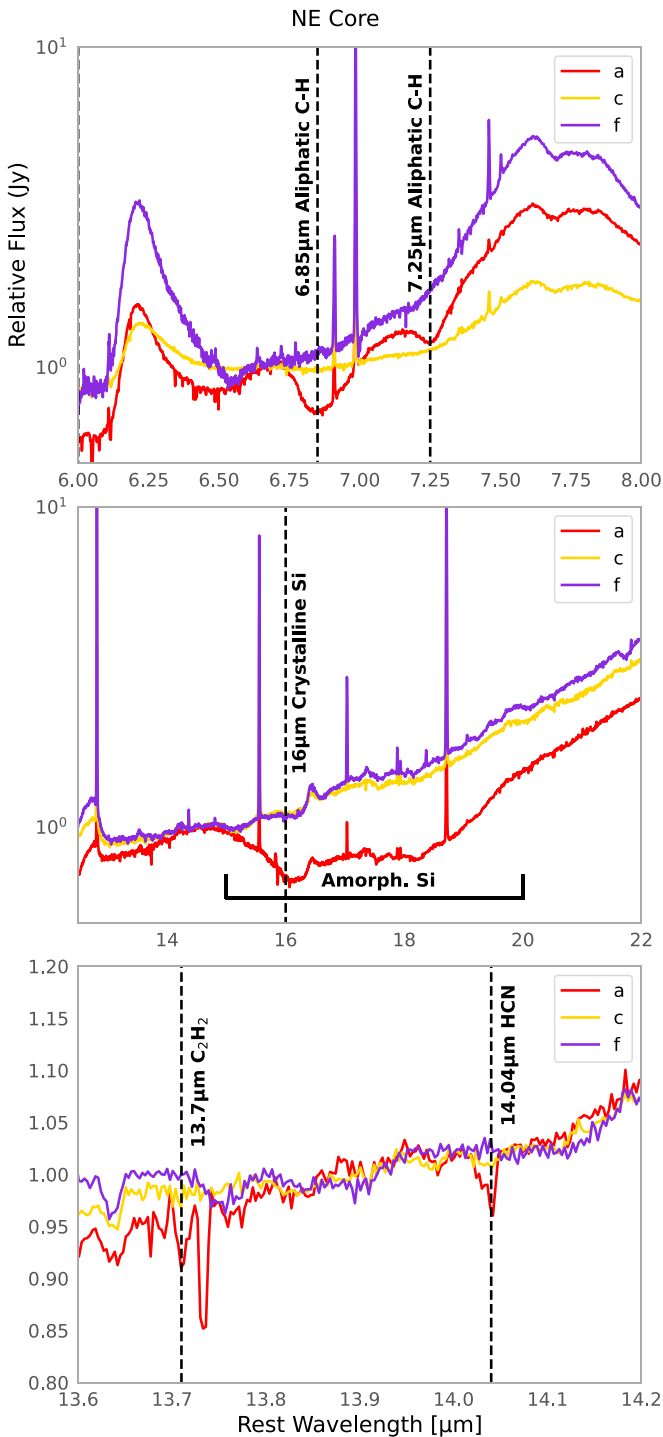


Figure 3. Expanded view of spectral features present only in the NE nucleus (region (a)) indicative of a highly embedded source, compared with spectra of the SW nucleus (c) and region (f) (embedded cluster). The dashed vertical lines in each panel correspond to the following features. Top panel: absorption features due to aliphatic hydrocarbons at 6.85 and 7.25 μm . The spectra are normalized to 1 Jy at 6.7 μm . Middle panel: crystalline silicate absorption features at 16 and 19 μm combined with amorphous silicate absorption at 18 μm . The spectra are normalized to 1 Jy at 14.5 μm . Bottom panel: C_2H_2 and HCN absorption. The spectra are normalized to 1 Jy at 13.8 μm .

recent ALMA study by Saito et al. (2018) does not detect HCN-vib emission toward the NE Core of VV 114 with a limit of $\Sigma(\text{HCN-vib}) < 0.12 L_{\odot} \text{pc}^{-2}$. The relatively faint millimeter continuum found by Saito et al. (2017) is consistent with the

HCN-vib non-detection. This suggests that the NE core either does not fulfill the CON criteria of Falstad et al. (2021) or that the radius of the CON region is smaller than 12 pc.

4. Discussion

The JWST spectra and imaging resolve a blend of diffuse emission from star formation and shocks, several reddened star-forming knots, and an AGN. Although the resolved spectra cover a much smaller region, our results confirm the analysis of MIRI imaging data of Evans et al. (2022). The SW core (c) has an elevated continuum at $\sim 5 \mu\text{m}$, consistent with the presence of dust in thermal equilibrium at temperatures near the sublimation temperature of silicate grains ($\sim 1200 \text{ K}$). This elevated continuum has been demonstrated by Laurent et al. (2000) and Petric et al. (2011) (and references therein) as a telltale sign of a dust enshrouded, optically thick AGN. We also find PAH emission in all of our extracted apertures and widespread diffuse emission in our emission line maps. The various emission and absorption features allow us to probe directly the nature of the bright cores and diffuse emission as well as the kinematics of the gas in VV 114E.

4.1. Evidence for AGN Activity

Previous multiwavelength observations of VV 114 hinted at an AGN contribution in the X-ray, mid-IR, and submillimeter (e.g., Le Floch et al. 2002; Grimes et al. 2006; Iono et al. 2013). As the brightest compact sources from the mid-IR to radio, the NE and SW cores are the most likely to harbor an AGN, with Iono et al. (2013) identifying the NE core as a potential AGN. We do not detect coronal lines in either the NE or SW cores, but this does not rule out an AGN; ULIRG Mrk 231 is a well-known optically classified Seyfert 1 with no observed coronal lines in the mid-IR. Instead, very low PAH EW and silicate absorption indicate the presence of the AGN in Mrk 231 (Armus et al. 2007; Imanishi et al. 2010; Inami et al. 2013; Stierwalt et al. 2014).

To compare with other AGN and starburst-dominated LIRGs we plot our measured values on several PAH diagnostic diagrams (Figure 4). We first plot $s_{9.7}$ versus $\text{EW}_{6.2\mu\text{m}}$ (Figure 4(d)): the apertures extracted from diffuse emission (with the exception of (d)) as well as the star-forming clump in aperture (f) have values consistent with starburst galaxies and lie near the literature values for M82 (Spoon et al. 2007). The SW core shows contribution from an AGN (e.g., Spoon et al. 2007; Marshall et al. 2018) and lies near Mrk 273, a ULIRG with a Seyfert 2 nucleus (Armus et al. 2007; Stierwalt et al. 2013; U et al. 2013). Aperture (d) is directly adjacent to the SW core, in fact partially overlapping, and is likely showing some contribution from the AGN that is more clearly seen in the SW core. The NE core has a correspondingly higher obscuration and lies near the values observed for ULIRG Arp 220, a deeply embedded starburst that also has C_2H_2 and HCN absorption crystalline silicate features in its mid-IR spectrum (Spoon et al. 2006; Lahuis et al. 2007).

The spectrum of the SW core at shorter IR wavelengths is consistent with the heating and processing of dust by an AGN, which reduces the 3.3 μm PAH EW (Figures 4(a)–(c)) as well as the 6.2 μm EW ($\text{EW}_{3.3\mu\text{m}} < 0.04 \mu\text{m}$ and $\text{EW}_{6.2\mu\text{m}} < 0.20 \mu\text{m}$; Imanishi et al. 2010; Petric et al. 2011). The 3.3 μm PAH feature is directly adjacent to a broad absorption feature caused by H_2O ice that may impact our

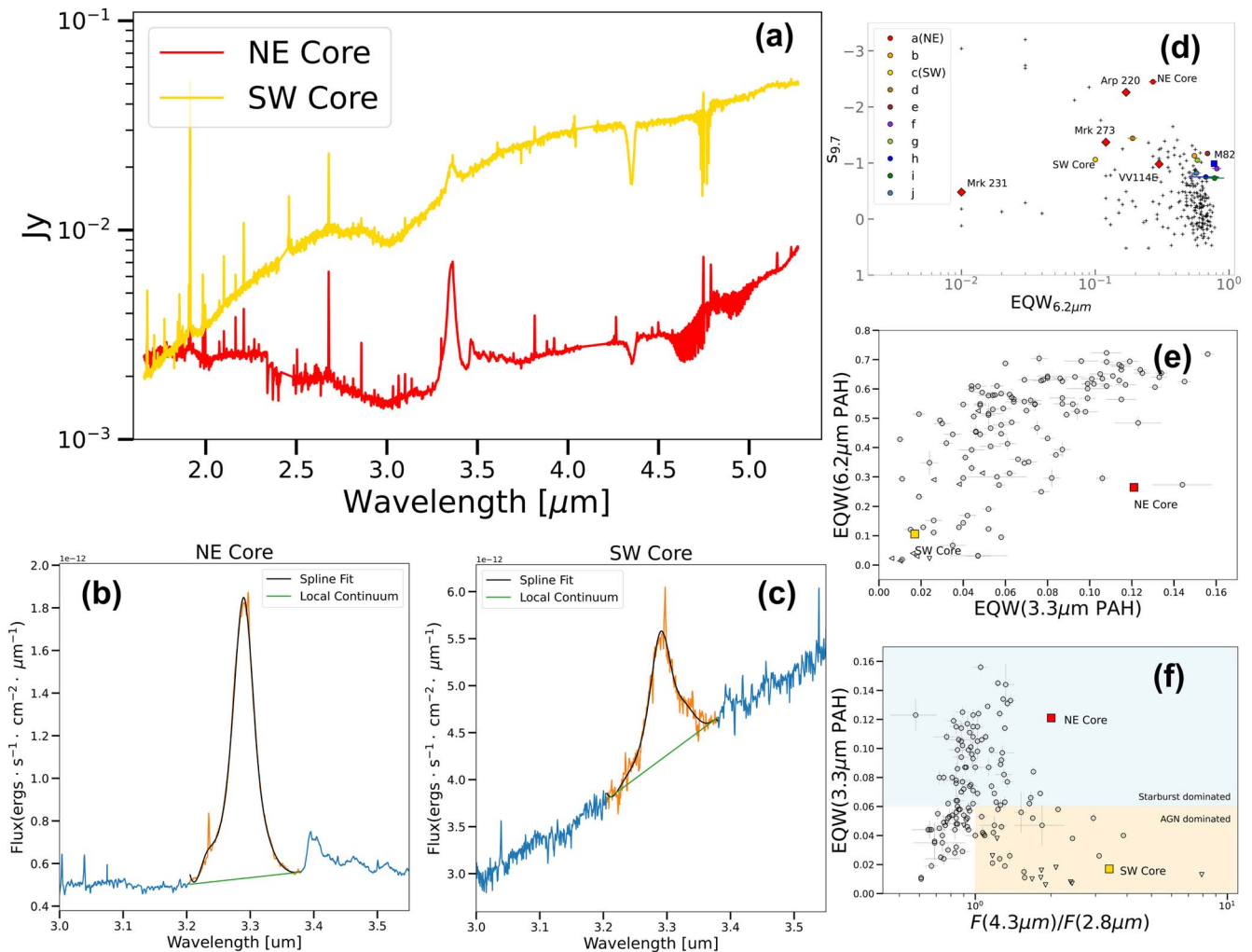


Figure 4. NIRSpect spectra and PAH diagnostic plots used. (a) NIRSpect spectra of the NE and SW cores in VV 114. The SW core shows a strongly rising continuum characteristic of hot dust associated with an AGN. The $3.3 \mu\text{m}$ PAH feature is detected in both spectra, but is significantly weaker in the SW core. (b)–(c) Expanded region and fit to the $3.3 \mu\text{m}$ PAH feature in the NE (b) and SW (c) cores. (d) EW of the $6.2 \mu\text{m}$ PAH feature plotted vs. silicate strength. Measurements for our MIRI apertures are plotted as colored circles. Values measured with Spitzer/IRS for ULIRGs and LIRGs from the GOALS sample are plotted as “+” symbols, with VV 114E and three comparative ULIRGs denoted with red diamonds (Stierwalt et al. 2013). The starburst galaxy M82 is shown as a blue square (Spoon et al. 2007). (e) EWs of the $6.2 \mu\text{m}$ and $3.3 \mu\text{m}$ PAH features for the SW and NE cores (Inami et al. 2018). (f) Flux density ratio and $3.3 \mu\text{m}$ PAH EW for the SW and NE cores (Inami et al. 2018). The SW core lies near AGN-dominated ULIRGs in both diagnostic plots.

continuum measurement and in turn the measured $\text{EW}_{3.3\mu\text{m}}$ (Imanishi et al. 2008; Inami et al. 2018). Moreover, the attenuation of the continuum and the $3.3 \mu\text{m}$ feature may be different depending on the geometry of the dust and PAH emission (Lai et al. 2020). We make a simple estimate of the impact by performing a power-law fit to the $1.7\text{--}5.0 \mu\text{m}$ spectrum and dividing our integrated flux by the continuum flux density at $3.3 \mu\text{m}$. This decreases the $\text{EW}_{3.3\mu\text{m}}$ slightly from $0.12 \mu\text{m}$ and $0.017 \mu\text{m}$ to $0.11 \mu\text{m}$ and $0.015 \mu\text{m}$ for the NE and SW core, respectively, but does not affect our conclusions regarding the nature the NE and SW cores.

Inami et al. (2018) used AKARI to propose revised AGN diagnostics including the $3.3 \mu\text{m}$ PAH feature as well as the $F_{\nu}(4.3)/F_{\nu}(2.8)$ flux density ratio. When placed on a plot of $\text{EW}_{6.2\mu\text{m}}$ versus $\text{EW}_{3.3\mu\text{m}}$ (Figure 4(e); following Inami et al. 2018), the SW core lies in a region populated by known AGNs while the high $\text{EW}_{3.3\mu\text{m}}$ of the NE core is more consistent with ULIRGs dominated by star formation. We also place the NE and SW cores on the $\text{EW}_{3.3\mu\text{m}}$ versus flux density ratio diagnostic proposed by Inami et al. (2018) in Figure 4(f). The

SW core again clearly lies in the AGN-dominated portion of the diagram.

To estimate the contribution of the AGN in the SW core to the total luminosity of VV 114 we take the estimation of $L_{\text{IR}} \sim 5 \pm 0.5 \times 10^{10} L_{\odot}$ by Evans et al. (2022), about 12% of the total L_{IR} , as an upper limit. The presence of PAH features in the SW core indicates a blended contribution to the IR of star formation and AGN. Sources with similar $\text{EW}_{3.3\mu\text{m}}$ and colors have a bolometric AGN contribution of $\sim 30\text{--}50\%$ (Díaz-Santos et al. 2017; Inami et al. 2018), which means the contribution of the AGN to the total luminosity of the entire VV 114 system is $\sim 5\%$.

Although Iono et al. (2013) identified the NE core as potentially harboring an AGN due to the enhanced HCN/HCO+ ratio, an analysis of X-ray and millimeter data by Privon et al. (2020) showed that in fact the HCN/HCO+ is not a robust indicator of total AGN luminosity or its fractional contribution to IR luminosity. Our findings regarding the nature of the two bright cores in VV 114E are not in agreement with the analysis of JWST MIRI data by Donnan et al. (2023) who

propose the NE core as a potential AGN host based on its compact nature.

Interestingly, the NE core does have a somewhat low $EW_{6.2\mu\text{m}}$ despite a strong $EW_{3.3\mu\text{m}}$, placing it in a region of Figure 4(e) with few other galaxies. The galaxies with the most similar values to the NE core are II Zw 96 and IRAS F19297-0406, which also have similar $s_{9.7\mu\text{m}}$ when comparing the values measured using Spitzer for all three galaxies (Stierwalt et al. 2013). II Zw 96 is an unusually compact and powerful starburst (Inami et al. 2010, 2022) and IRAS F19297-0406 has a powerful starburst-driven outflow (Soto et al. 2012; Veilleux et al. 2013), physical conditions which may warm dust in a way that lowers $EW_{6.2\mu\text{m}}$. The curious nature of both the NE and SW cores warrants a more thorough follow-up analysis decomposing the relative contributions of the stellar, PAH, dust, and AGN components of the SED.

4.2. Shocks and Tidal Features

Evidence for extended shocks in VV 114 has previously been suggested at visible wavelengths via enhanced [O I] and [S II] emission and broadened line profiles across VV 114 (Rich et al. 2011, 2015) and in the submillimeter from the presence of methanol (CH_3OH) in the “overlap” region between VV 114W and E. Apertures (f), (g), (h), and (i) lie closest to the overlap region and both emission line ratios and molecular and atomic line widths show evidence of shocked gas.

The atomic lines in aperture (f) all have narrow line widths ($\sim 100\text{--}150\text{ km s}^{-1}$) and ratios typical of buried, young star-forming regions as this source was revealed to be in (Linden et al. 2022). Aperture (i) encompasses a fainter star-forming knot seen in both MIRI and NIRCAM imaging, and several emission line profiles show a double peaked profile. The broader H_2 lines in apertures (f) and (i) ($\sim 200\text{--}300\text{ km s}^{-1}$) are similar to the values extracted just to the north in apertures (g) and (h), which show elevated [Fe II]/ $\text{Pf}\alpha$ values indicative of shocks, also seen in MRS observations of NGC 7469 (U et al. 2022).

These regions are consistent with ALMA observations of VV 114 where the elevated FWHM of $150\text{--}300\text{ km s}^{-1}$ in the molecular gas (e.g., $\text{CO}(1\text{--}0)$, CH_3OH) is found in the “overlap” region and is suggested to be the product of both shocks, and overlapping kinematic components due to the merger (Saito et al. 2017, 2018) and may be similar in nature to the shocked bridge seen in Stephan’s Quintet (Guillard et al. 2012b; Appleton et al. 2017). However, the clusters in this region of VV 114 are 1–2 orders of magnitude more massive and $\sim 2\text{--}3$ times dustier than those seen in the bridge of Stephan’s Quintet (Fedotov et al. 2011; Linden et al. 2022).

Resolved emission line profiles in aperture (j) show a strong indication of both blue and redshifted wings, potentially due to projection effects of the tidal arm that extends from VV 114W across the IR-bright cores southward beyond the FOV of our aperture extractions (Saito et al. 2015). These values are supported by examining the [Ne II] emission line profiles which fall in Channel 3 and has a wider FOV than Channel 1. Looking several arcseconds southeast of the SW core, the [Ne II] line profiles show significant broadening and wings in several spaxels (Figure 2). We also see enhanced H_2 /[Ne II] in this region and at the furthest SE region of our spaxel-by-spaxel maps. The value of [Fe II]/ $\text{Pf}\alpha \sim 9$ in aperture (j) also indicates the presence of shocked gas which likely extends

beyond the Channel 1 FOV and follows the elevated emission line ratios caused by shocks seen in visible light IFU observations (Rich et al. 2011).

Aperture (b) appears to be dominated by diffuse emission when examining the MIRI and NIRCAM images, but does not display the same characteristics of shock excitation as the other apertures that trace diffuse gas. Previous observations of shock excitation in VV 114 have suggested both galactic winds and tidally driven gas flows as sources of shock excitation (Rich et al. 2011; Saito et al. 2017). The JWST data show some shocked gas coincident with tidal features previously observed in the submillimeter and radio as well as kinematic profiles potentially associated with galactic winds. Follow-up work mapping the 2D kinematics of the molecular and atomic gas, as well as the temperature and distribution of H_2 gas using these data will provide a more complete picture of the shock excitation in VV 114, especially when combined with wide field optical IFU observations from MUSE.

4.3. SFRs

If we assume that the atomic and fine structure line emission is dominated by star formation in the apertures with bright star-forming clumps ((a), (e), and (f)), we can calculate star formation rates (SFRs) using a hydrogen recombination line flux or neon emission. If we scale either Pfund or Humphreys α to $\text{H}\alpha$ assuming Case B recombination (Hummer & Storey 1987) and use Equation (2) of Murphy et al. (2011) the estimated SFR for each spectrum is $\sim 1 M_\odot \text{ yr}^{-1}$. Using the [Ne II] and [Ne III] fluxes and Equation (3) of Ho & Keto (2007) yields an SFR from $\sim 1.5\text{--}2.5 M_\odot \text{ yr}^{-1}$. These values are consistent with those found by Song et al. (2022) using the radio emission of bright knots measured with the VLA and amount to $\sim 2\%\text{--}3\%$ of the total SFR per aperture.

5. Conclusions

Our analysis of MIRI MRS and NIRSpect IFU spectroscopy of VV 114E shows emission and absorption features that allow us to resolve variations in the properties of the IR-bright nuclear cores, unresolved clumps, and diffuse gas. The integrated properties of the two bright cores and the diffuse gas agree with past multiwavelength observations of VV 114E that show widespread star formation and diffuse emission from shocks, and reveal spectroscopic evidence of an AGN in the SW core. More specifically:

1. The SW core harbors an AGN as indicated by its extremely low 3.3 and $6.2\mu\text{m}$ PAH EWs and strong $3\text{--}5\mu\text{m}$ continuum, consistent with AGN-dominated LIRGs. The SW core is also surrounded by star-forming knots and diffuse emission, which are evident in the atomic line ratios at longer wavelengths. The AGN in the SW core likely accounts for $\sim 5\%$ of the total luminosity of VV 114.
2. The NE core is deeply embedded, its mid-IR spectrum displays strong $9.7\mu\text{m}$ silicate absorption, crystalline silicate features, aliphatic hydrocarbons, and HCN absorption. Using the $14\mu\text{m}$ HCN absorption line we calculate a temperature and column density of $300\text{--}500\text{ K}$ and $N(\text{HCN})$ of $1\text{--}5 \times 10^{16}\text{ cm}^{-2}$, respectively. Our data show no evidence of an AGN impacting the atomic or molecular gas at mid-IR wavelengths and we conclude that this source is a deeply buried star-forming region.

3. The diffuse gas NW of the nuclear region shows elevated [Fe II]/Pfund α and higher H₂ line widths along with double peaked profiles caused by shocked gas in the overlap region previously observed by ALMA. A fit to [Ne II] across the MIRI FOV reveals broader emission profiles to the south of the nucleus consistent with the extended tidal feature observed by ALMA, as well as shocked gas $\sim 1''$ SE of the SW core that likely extends beyond the Channel 1 FOV, consistent with previous visible light IFU observations.

These early 3D spectral data highlight the power of combining the NIRSpect and MIRI data to elucidate the nature of complex, obscured star formation and AGNs in the local universe. Taken together the spectroscopic data sets from both JWST instruments are extremely rich and will facilitate detailed and thorough analysis in future papers in this series.

We thank the referee for their helpful comments. This work is based on observations made with the NASA/ESA/CSA JWST. The research was supported by NASA grant JWST-ERS-01328. The data were obtained from the Mikulski Archive for Space Telescopes at the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-03127 for JWST. The specific observations analysed can be accessed via [10.17909/yqk1-jr92](https://doi.org/10.17909/yqk1-jr92). V.U. acknowledges funding support from NASA Astrophysics Data Analysis Program (ADAP) grant 80NSSC20K0450. The Flatiron Institute is supported by the Simons Foundation. A.M.M. acknowledges support from the National Science Foundation under grant No. 2009416. A.S.E. and S.L. acknowledge support from NASA grants HST-GO15472 and HST-GO16914. Y.S. was funded in part by the NSF through the Grote Reber Fellowship Program administered by Associated Universities, Inc./National Radio Astronomy Observatory. The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc. F.M.-S. acknowledges support from NASA through ADAP award 80NSSC19K1096. S.A. gratefully acknowledges support from an ERC Advanced grant 789410, from the Swedish Research Council and from the Knut and Alice Wallenberg (KAW) Foundation. S.A. gratefully acknowledges John Black for helpful discussions. K.I. acknowledges support by the Spanish MCIN under grant PID2019-105510GB-C33/AEI/10.13039/501100011033. This work was also partly supported by the Spanish program Unidad de Excelencia Mar de Maeztu CEX2020-001058-M, financed by MCIN/AEI/10.13039/501100011033. The computations presented here were conducted through Carnegie's partnership in the Resnick High Performance Computing Center, a facility supported by Resnick Sustainability Institute at the California Institute of Technology. Finally, this research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.

Facilities: JWST (NIRCam, NIRSpect and MIRI).

ORCID iDs

J. Rich <https://orcid.org/0000-0002-5807-5078>
S. Aalto <https://orcid.org/0000-0002-5828-7660>

A. S. Evans <https://orcid.org/0000-0003-2638-1334>
V. Charmandaris <https://orcid.org/0000-0002-2688-1956>
G. C. Privon <https://orcid.org/0000-0003-3474-1125>
T. Lai <https://orcid.org/0000-0001-8490-6632>
H. Inami <https://orcid.org/0000-0003-4268-0393>
S. Linden <https://orcid.org/0000-0002-1000-6081>
L. Armus <https://orcid.org/0000-0003-3498-2973>
T. Diaz-Santos <https://orcid.org/0000-0003-0699-6083>
P. Appleton <https://orcid.org/0000-0002-7607-8766>
L. Barcos-Muñoz <https://orcid.org/0000-0003-0057-8892>
T. Böker <https://orcid.org/0000-0002-5666-7782>
K. L. Larson <https://orcid.org/0000-0003-3917-6460>
D. R. Law <https://orcid.org/0000-0002-9402-186X>
M. A. Malkan <https://orcid.org/0000-0001-6919-1237>
A. M. Medling <https://orcid.org/0000-0001-7421-2944>
Y. Song <https://orcid.org/0000-0002-3139-3041>
V. U <https://orcid.org/0000-0002-1912-0024>
P. van der Werf <https://orcid.org/0000-0001-5434-5942>
M. J. I. Brown <https://orcid.org/0000-0002-1207-9137>
L. Finnerty <https://orcid.org/0000-0002-1392-0768>
C. Hayward <https://orcid.org/0000-0003-4073-3236>
J. Howell <https://orcid.org/0000-0001-6028-8059>
K. Iwasawa <https://orcid.org/0000-0002-4923-3281>
F. Kemper <https://orcid.org/0000-0003-2743-8240>
J. Marshall <https://orcid.org/0000-0001-7712-8465>
J. M. Mazzarella <https://orcid.org/0000-0002-8204-8619>
J. McKinney <https://orcid.org/0000-0002-6149-8178>
F. Muller-Sanchez <https://orcid.org/0000-0002-2713-0628>
E. J. Murphy <https://orcid.org/0000-0001-7089-7325>
D. Sanders <https://orcid.org/0000-0002-1233-9998>
S. Stierwalt <https://orcid.org/0000-0002-2596-8531>
J. Surace <https://orcid.org/0000-0001-7291-0087>

References

- Aalto, S., Martín, S., Costagliola, F., et al. 2015, *A&A*, 584, A42
Aalto, S., Muller, S., König, S., et al. 2019, *A&A*, 627, A147
Appleton, P. N., Guillard, P., Togi, A., et al. 2017, *ApJ*, 836, 76
Armus, L., Charmandaris, V., Bernard-Salas, J., et al. 2007, *ApJ*, 656, 148
Armus, L., Lai, T., V., U., et al. 2023, *ApJL*, 942, L37
Arp, H. 1966, *ApJS*, 14, 1
Böker, T., Arribas, S., Lützgendorf, N., et al. 2022, *A&A*, 661, A82
Bouchet, P., García-Marín, M., Lagage, P.-O., et al. 2015, *PASP*, 127, 612
Bushouse, H., Eisenhamer, J., Dencheva, N., et al. 2022, *spacetelescope/jwst*: JWST v1.6.2, Zenodo, doi:10.5281/zenodo.6984366
Charmandaris, V., Le Floch, E., & Mirabel, I. F. 2004, *ApJL*, 600, L15
Díaz-Santos, T., Armus, L., Charmandaris, V., et al. 2017, *ApJ*, 846, 32
Donnan, F. R., García-Berete, I., Rigopoulou, D., et al. 2023, *MNRAS*, 519, 3691
Evans, A. S., Frayer, D., Charmandaris, V., et al. 2022, *ApJL*, 940, L8
Falstad, N., Aalto, S., König, S., et al. 2021, *A&A*, 649, A105
Fedotov, K., Gallagher, S. C., Konstantopoulos, I. S., et al. 2011, *AJ*, 142, 42
Garofali, K., Lehmer, B. D., Basu-Zych, A., et al. 2020, *ApJ*, 903, 79
Genzel, R., Lutz, D., Sturm, E., et al. 1998, *ApJ*, 498, 579
Grimes, J. P., Heckman, T., Hoopes, C., et al. 2006, *ApJ*, 648, 310
Guillard, P., Ogle, P. M., Emonts, B. H. C., et al. 2012a, *ApJ*, 747, 95
Guillard, P., Boulanger, F., Pineau des Forêts, G., et al. 2012b, *ApJ*, 749, 158
Harada, N., Thompson, T. A., & Herbst, E. 2013, *ApJ*, 765, 108
Ho, L. C., & Keto, E. 2007, *ApJ*, 658, 314
Howell, J. H., Armus, L., Mazzarella, J. M., et al. 2010, *ApJ*, 715, 572
Imanishi, M., Nakagawa, T., Ohya, Y., et al. 2008, *PASJ*, 60, S489
Imanishi, M., Nakagawa, T., Shirahata, M., Ohya, Y., & Onaka, T. 2010, *ApJ*, 721, 1233
Inami, H., Armus, L., Surace, J. A., et al. 2010, *AJ*, 140, 63
Inami, H., Armus, L., Charmandaris, V., et al. 2013, *ApJ*, 777, 156
Inami, H., Armus, L., Matsuhara, H., et al. 2018, *A&A*, 617, A130
Inami, H., Surace, J., Armus, L., et al. 2022, *ApJL*, 940, L6
Iono, D., Saito, T., Yun, M. S., et al. 2013, *PASJ*, 65, L7
Jakobsen, P., Ferruit, P., Alves de Oliveira, C., et al. 2022, *A&A*, 661, A80

- Labiano, A., Argyriou, I., Álvarez-Márquez, J., et al. 2021, *A&A*, **656**, A57
- Lahuis, F., & van Dishoeck, E. F. 2000, *A&A*, **355**, 699
- Lahuis, F., Spoon, H. W. W., Tielens, A. G. G. M., et al. 2007, *ApJ*, **659**, 296
- Lai, T. S.-Y., Smith, J. D. T., Baba, S., Spoon, H. W. W., & Imanishi, M. 2020, *ApJ*, **905**, 55
- Laurent, O., Mirabel, I. F., Charmandaris, V., et al. 2000, *A&A*, **359**, 887
- Le Floc'h, E., Charmandaris, V., Laurent, O., et al. 2002, *A&A*, **391**, 417
- Linden, S. T., Evans, A. S., Armus, L., et al. 2022, arXiv:2210.05763
- Lutz, D., Sturm, E., Genzel, R., et al. 2000, *ApJ*, **536**, 697
- Marshall, J. A., Elitzur, M., Armus, L., Diaz-Santos, T., & Charmandaris, V. 2018, *ApJ*, **858**, 59
- Murphy, E. J., Condon, J. J., Schinnerer, E., et al. 2011, *ApJ*, **737**, 67
- Newville, M., Stensitzki, T., Allen, D. B., & Ingargiola, A. 2014, LMFIT: Non-Linear Least-Square Minimization and Curve-Fitting for Python, 0.8.0, Zenodo, doi:10.5281/zenodo.11813
- Petric, A. O., Armus, L., Howell, J., et al. 2011, *ApJ*, **730**, 28
- Privon, G. C., Ricci, C., Aalto, S., et al. 2020, *ApJ*, **893**, 149
- Ricci, C., Privon, G. C., Pfeifle, R. W., et al. 2021, *MNRAS*, **506**, 5935
- Rich, J. A., Kewley, L. J., & Dopita, M. A. 2011, *ApJ*, **734**, 87
- Rich, J. A., Kewley, L. J., & Dopita, M. A. 2015, *ApJS*, **221**, 28
- Rieke, G. H., Wright, G. S., Boker, T., et al. 2015, *PASP*, **127**, 584
- Saito, T., Iono, D., Yun, M. S., et al. 2015, *ApJ*, **803**, 60
- Saito, T., Iono, D., Espada, D., et al. 2017, *ApJ*, **834**, 6
- Saito, T., Iono, D., Espada, D., et al. 2018, *ApJ*, **863**, 129
- Sakamoto, K., Aalto, S., Evans, A. S., Wiedner, M. C., & Wilner, D. J. 2010, *ApJL*, **725**, L228
- Sakamoto, K., Gonzalez-Alfonso, E., Martin, S., et al. 2021, *ApJ*, **923**, 206
- Schilke, P., Walmsley, C. M., Pineau Des Forets, G., et al. 1992, *A&A*, **256**, 595
- Soifer, B. T., Sanders, D. B., Madore, B. F., et al. 1987, *ApJ*, **320**, 238
- Song, Y., Linden, S. T., Evans, A. S., et al. 2022, *ApJ*, **940**, 52
- Soto, K. T., Martin, C. L., Prescott, M. K. M., & Armus, L. 2012, *ApJ*, **757**, 86
- Spoon, H. W. W., Marshall, J. A., Houck, J. R., et al. 2007, *ApJL*, **654**, L49
- Spoon, H. W. W., Tielens, A. G. G. M., Armus, L., et al. 2006, *ApJ*, **638**, 759
- Stierwalt, S., Armus, L., Surace, J. A., et al. 2013, *ApJS*, **206**, 1
- Stierwalt, S., Armus, L., Charmandaris, V., et al. 2014, *ApJ*, **790**, 124
- Sturm, E., Lutz, D., Verma, A., et al. 2002, *A&A*, **393**, 821
- U, V., Medling, A., Sanders, D., et al. 2013, *ApJ*, **775**, 115
- U, V., Lai, T., Bianchin, M., et al. 2022, *ApJ*, **940**, L5
- Veilleux, S., Melendez, M., Sturm, E., et al. 2013, *ApJ*, **776**, 27
- Weedman, D. W., Hao, L., Higdon, S. J. U., et al. 2005, *ApJ*, **633**, 706
- Wells, M., Pel, J.-W., Glasse, A., et al. 2015, *PASP*, **127**, 646
- Wright, E. L. 2006, *PASP*, **118**, 1711